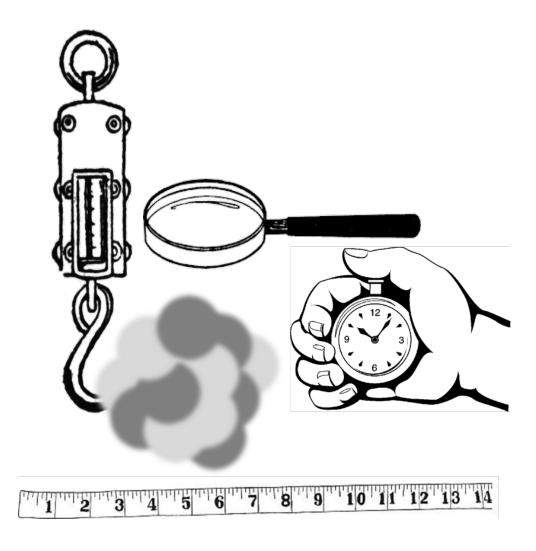
Lecture 1: *Nuclear Properties*

- General descriptors of nuclei
- Nuclear phenomenology
- The nuclear landscape



The Atomic Nucleus

- A conglomeration of *nucleons*, a.k.a. neutrons and protons
- Convention for referring to a particular configuration, a.k.a *nuclide*
 - $_{Z}^{A}(chemical\ symbol)_{N}$
 - A = # of nucleons, Z = # of protons, N = # of neutrons.
 - The chemical symbol is from the periodic table that corresponds to Z
 - Since A=Z+N and Z is indicated by the chemical symbol, more common notation is: ^Achemical symbol
 - For example, the most common type of carbon (Z=6) has 6 neutrons: ¹²C Conversationally, you would call this "carbon twelve"
- Nuclides with the same Z but different N are isotopes (though this term is often used in lieu of nuclides)
- Nuclides with the same N but different Z are isotones
- Nuclides with the same A are isobars
- Nuclides with nicknames:
 - ${}^{1}H = proton$, ${}^{2}H = deuteron (d)$, ${}^{3}H = triton (t)$, ${}^{4}He = \alpha$, ${}^{3}He = helion (rarely used)$

An "m" after the mass-number of a nucleus indicates an isomer, a long-lived excited state. E.g. 26mAl

How big is a nucleus?

The definition of u being based on ¹²C means it is a valuable tool for high-precision mass measurements, e.g. C.Scheidenberger et al., Nuc. Phys. A 2002

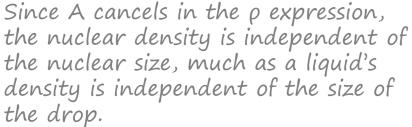
Phenomenological estimates:

- A nucleus's mass is roughly: M(Z,A) = A*amu
 - 1amu = atomic mass unit = u = $931.494 \text{MeV/c}^2 \approx 1.66 \times 10^{-24} \text{g}$
 - The amu is defined such that $M(^{12}C) \equiv 12u$
- A nucleus's (charge) radius is roughly: $R(Z,A) = (1.2fm)*A^{1/3}$
 - fm = femtometer (a.k.a. fermi) = 10^{-15} m
 - The radius of a nucleon is often referred to as $r_0=1.2$ fm
 - For the RMS radius, multiply by $\sqrt{3/5}$
- Therefore, an estimate for the nuclear density is:

•
$$\rho = \frac{M}{V} = \frac{M}{\frac{4}{3}\pi R^3} = \frac{(1u)A}{\frac{4}{3}\pi (1.2fm)^3 A} \approx 0.14 \text{ nucleons/fm}^3$$

• For fun, in terms of mass-density: $\rho = \frac{(1.66 \times 10^{-24} g)A}{\frac{4}{3}\pi (1.2 fm)^3 A} \approx 2.3 \times \frac{10^{-25} g}{fm^3}$

...which doesn't sound like much, but this is $2x10^{14}$ g/cm³ (the Great Pyramid of Giza is only ~10¹² grams)



Partly inspired by this property, some basic nuclear calculations are based on this liquid drop analogy. (G. Gamow, Proc. Roy. Soc. A 1929,1930)

Nuclear Transmutation

- Rules for converting one nuclide (or nuclides) to another (or others)
 - Charge conservation: $\sum q_{before} = \sum q_{after}$ (q from protons + positrons + electrons)
 - Baryon conservation: $\sum A_{before} = \sum A_{after}$ (A from neutrons + protons)
 - Lepton number conservation: $[N_{leptons} N_{anti-leptons}]_{before} = [N_{leptons} N_{anti-leptons}]_{after}$
 - * Transmutation likelihoods are impacted by energetics and spin/parity selection rules

Two Types:

- Reactions
 - Multiple reactants create one or more products
 - Notation: A+b \rightarrow c+D is written as A(b,c)D, where M(b)<M(A) and M(c)<M(D)
 - E.g. $^{12}\text{C}+\alpha \rightarrow ^{16}\text{O}+\gamma$ is $^{12}\text{C}(\alpha,\gamma)^{16}\text{O}$ or even just $^{12}\text{C}(\alpha,\gamma)$ and is called "carbon-twelve alpha gamma"
- Decays
 - α , β^+ , β^- , e^- -capture, β -delayed $\gamma/p/\alpha/n$ emission, fission, cluster emission, prompt γ
 - Lose nucleons for all above except β decay, e⁻-capture, and prompt γ (following a reaction)

The first nuclear reaction intentionally made in the laboratory was $^{14}N(\alpha,p)$ in 1919. (E. Rutherford, Nature 1935).

The first measured radioactive decay was α decay from uranium. (H. Becquerel, Comptes Rendus 1896).

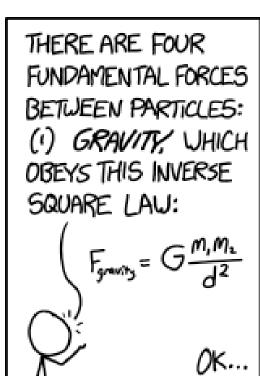
Nuclear Forces, mechanisms for binding & transmutation

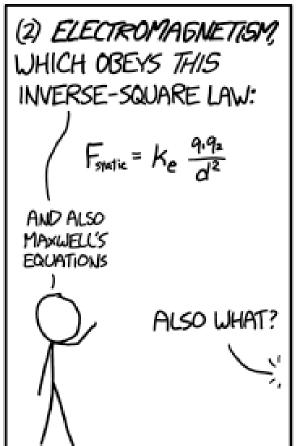
Fundamental forces of nature:

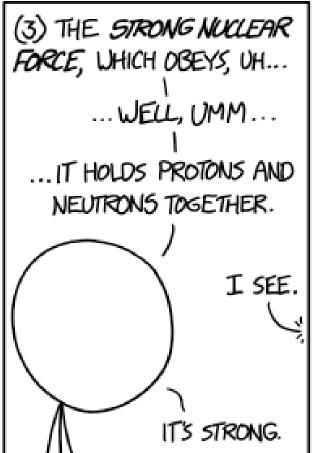
Force	Range (m)	Relative strength	Force carrier, X
Gravitational	∞	10 ⁻³⁸	Graviton
Weak	10 ⁻¹⁸	10 ⁻⁵	W [±] , Z ⁰
Electromagnetic	∞	$\alpha \approx 1/137$	Photon
Strong	10 ⁻¹⁵	1	Gluon, Pion

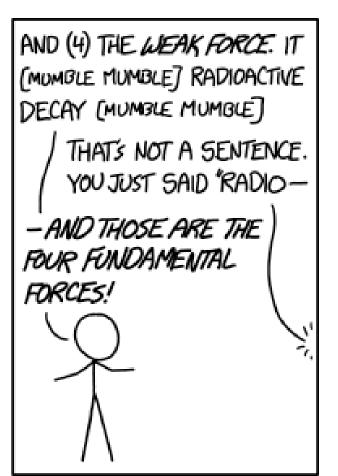
The range comes from the uncertainty principle:

- When converting reactant(s) to product(s), energy is spent on the mass of the force carrier: $\Delta E \ge M_\chi c^2$
- Violating energy conservation like this is only allowed due to the uncertainty principle: $\Delta E \Delta t \ge \hbar/2$
- And therefore the force carrier must be reabsorbed before: $t \le \hbar/(2\Delta E)$
- Since a particle can't be moving faster than the speed of light, c, the furthest range a force can be mediated is: $r \le tc = \hbar c/(2\Delta E) = \hbar c/(2M_X c^2) = \hbar/(2M_X c)$
- Thus, since factors of 2 are for chumps, the range of an interaction is defined as: $R \equiv \hbar/(M_x c)$
- Since the graviton and photon are massless, $R = \infty$. The trick for these calculations is that $\hbar c \approx 197 \text{MeV*fm}$.
- For the weak force, $M_X \sim 100 \text{GeV/c}^2$, so $R_{\text{weak}} \sim 10^{-3} \text{ fm}$.
- Similarly, since $M_{\pi} \sim 100 \text{MeV/c}^2$, $R_{\text{strong}} \sim 1 \text{fm}$.

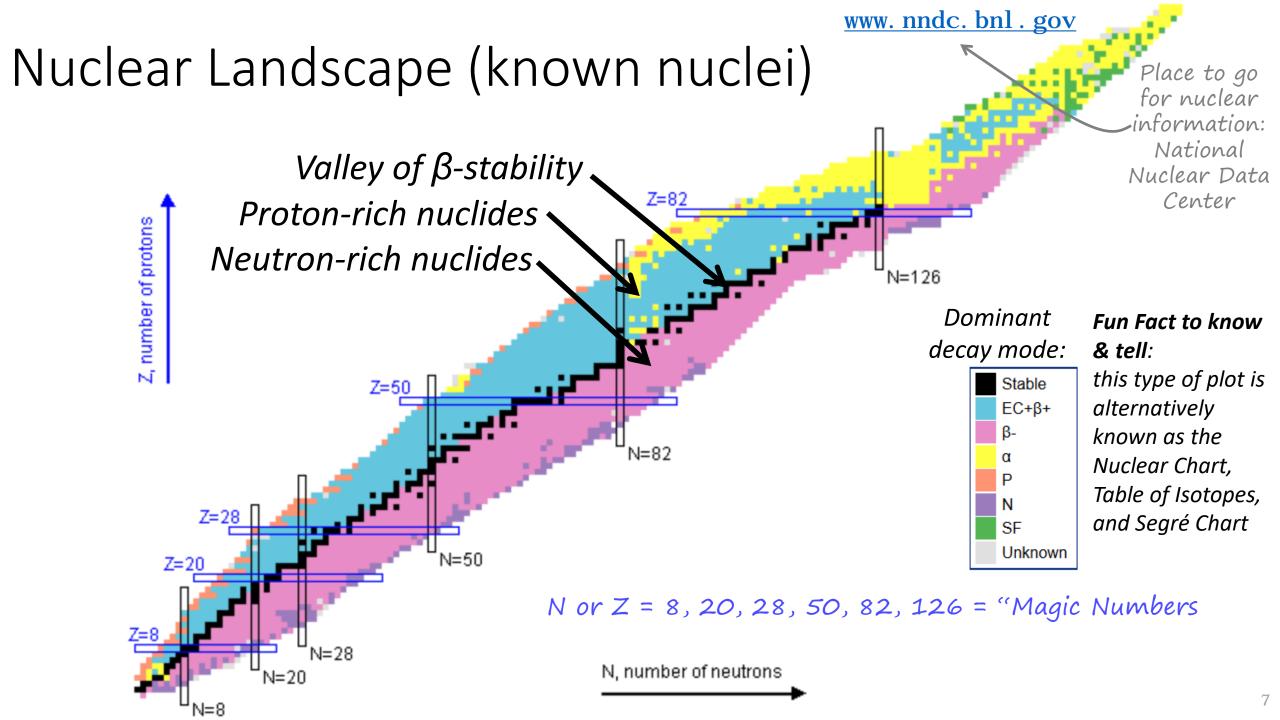


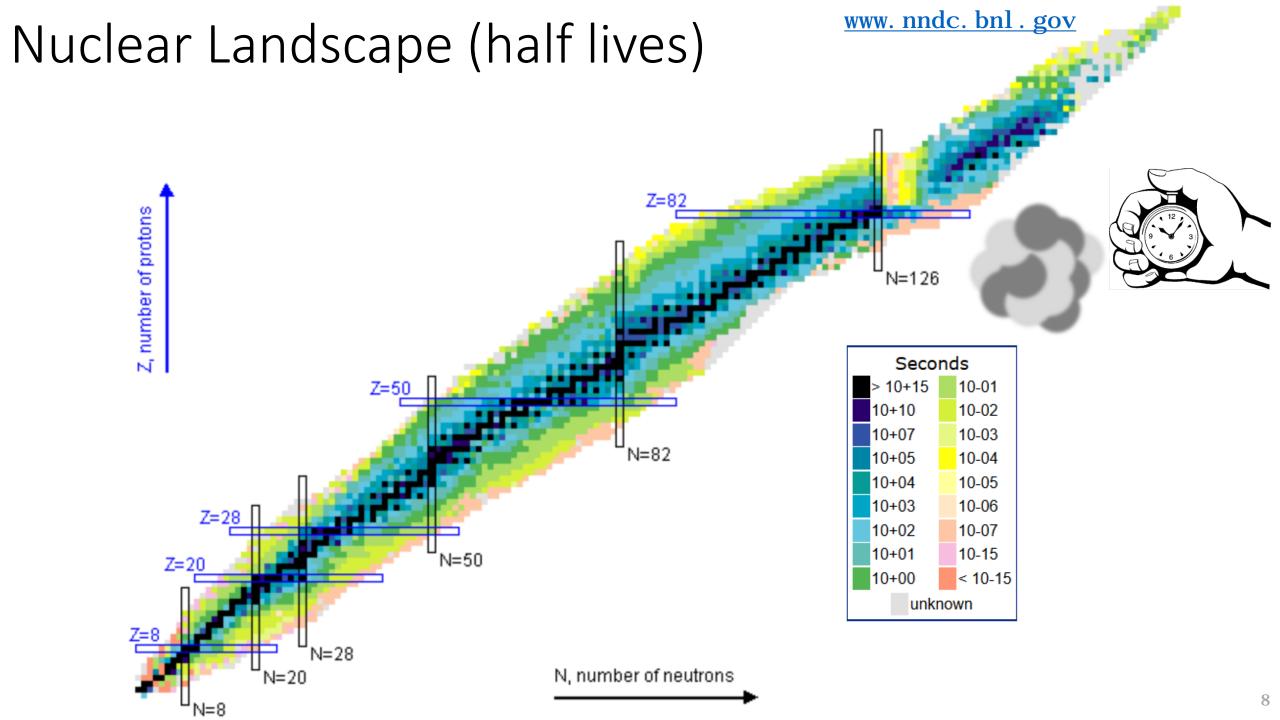




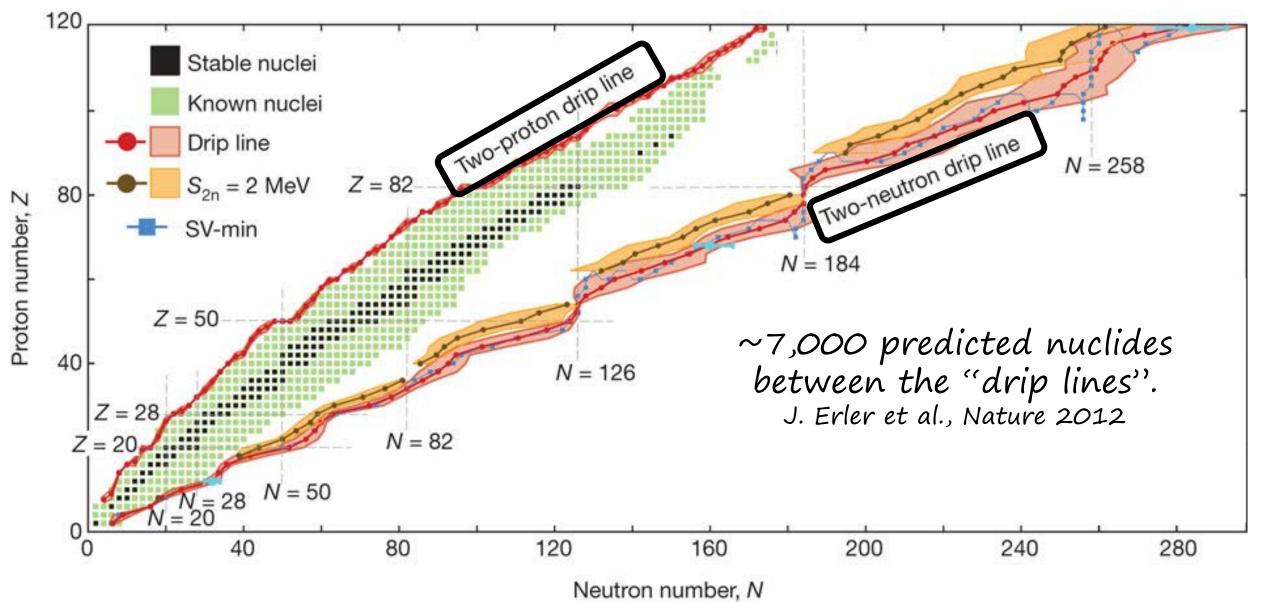


Randall Munroe, xkcd. com



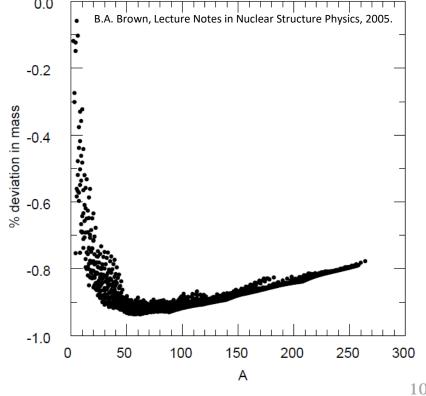


Nuclear Landscape (predicted nuclei)



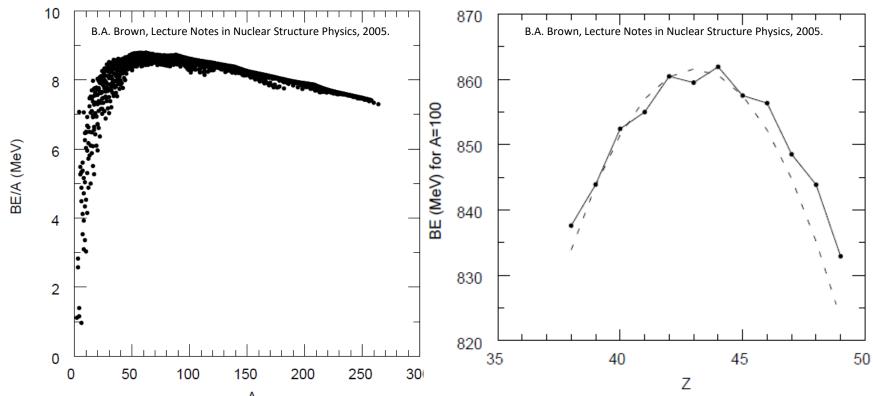
Nuclear Binding Energy

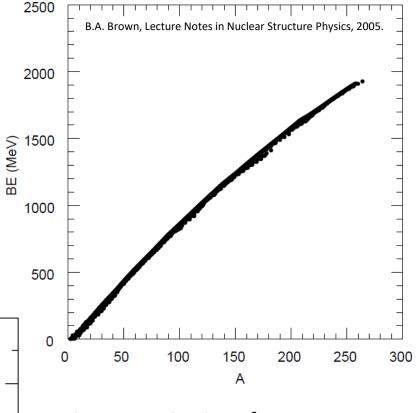
- Early mass spectrometry showed that nuclear masses are (nearly) integer multiples of the hydrogen atom mass. (Coined the "whole number rule" (F. Aston, Nature 1919))
- However, high-precision work showed deviations on the sub-% level (F. Aston Nature 1924)
- This deviation is known as the "mass defect", whereby the nuclear mass is a bit less than the sum of the constituent nucleon masses:
- Which is of course a result of Einstein's postulate, $E=mc^2$
- Nucleons within a nucleus are bound together, and this binding requires energy.
- The binding energy, BE, is paid-for via reduced mass.
- Indeed, the lower-energy state of a nucleus is the only reason nucleons cluster together at all.



Nuclear Binding Energy

- The binding energy increases nearly linearly with A:
- It's not a surprise that more nucleons would lead to more binding, since nucleons will be prone to attraction via the strong force.
- As such, a more interesting quantity is BE/A



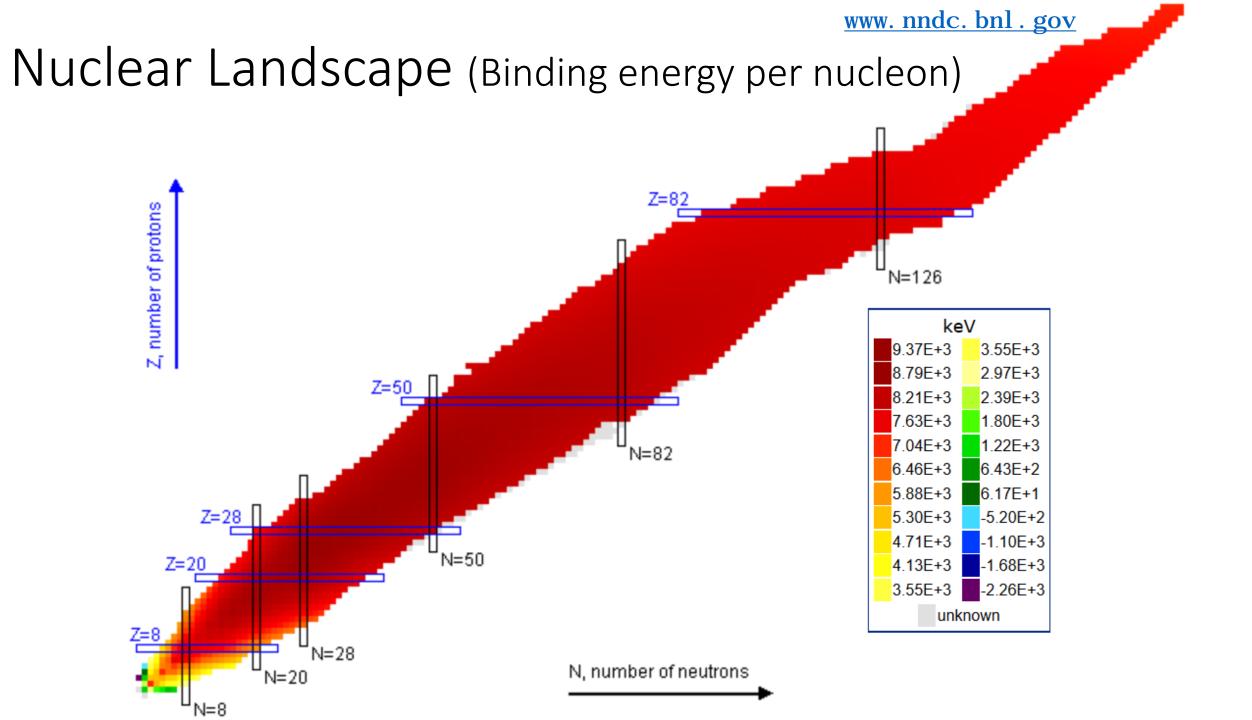


What are the key features you notice on the two left plots?

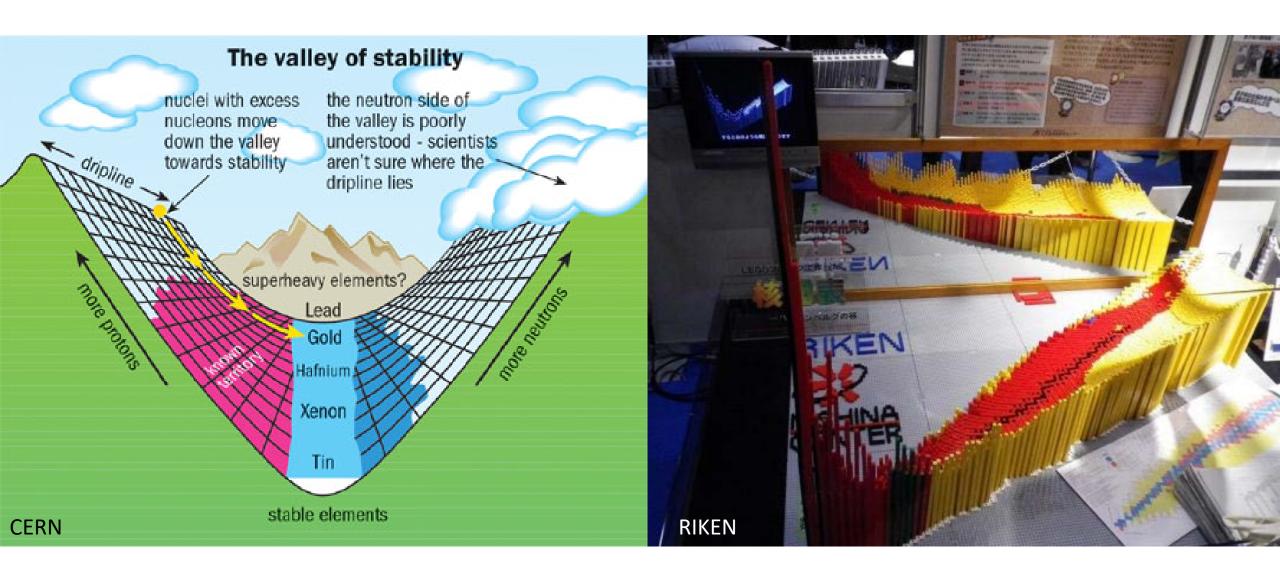
What to they imply about

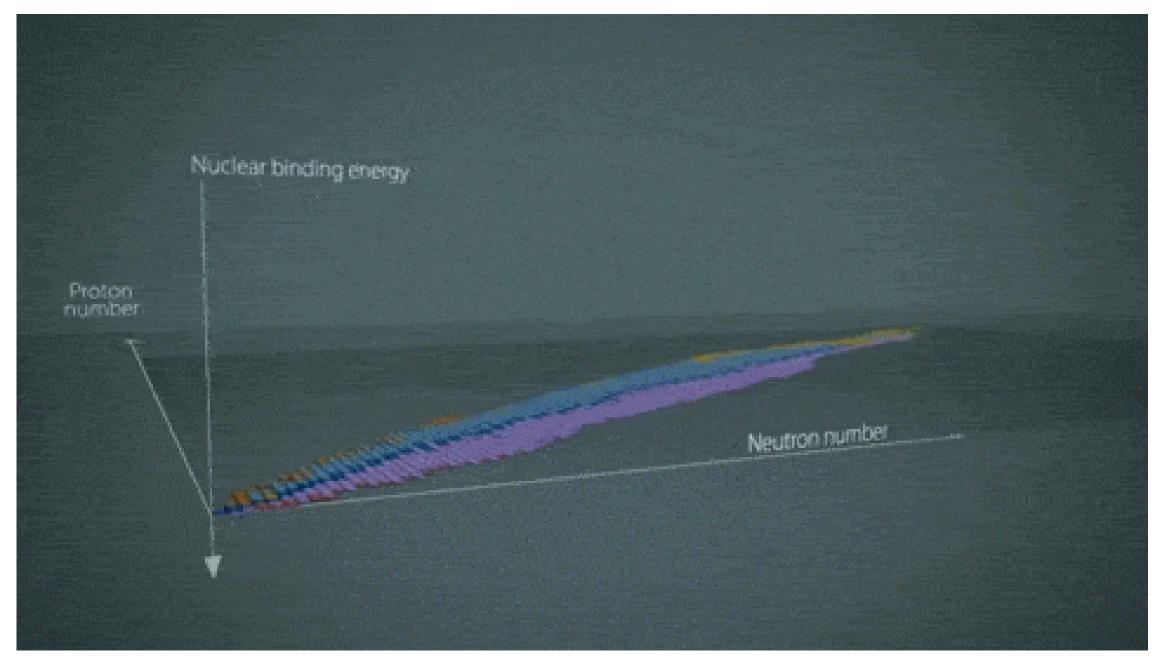
- · stability to decay?
- fusion?

Why do nuclei other than the peak BE/A form at all?



Binding Energy per nucleon: sort of a valley or half-pipe





F. Legrand & F. Durillon, https://www.youtube.com/watch?v=UTOp 2ZVZmM

Aside on nuclear masses, atomic mass excess

- Rather than binding energy, "atomic mass excess" is more commonly used.
- The conversion can be accomplished with the following equations:
 - 1. $BE(Z,A) \equiv Z^*(m_p + m_e) + N^*m_n M(Z,N)$
 - 2. $ME(Z,A) \equiv M(Z,N) (Z+N)*m_u$

here M(Z,N) is the mass of a nucleus with Z protons and N neutrons,

 $m_p = 938.272000 \text{ MeV/c}^2$ is the proton mass

 $m_e = 0.51100 \text{ MeV/c}^2$ is the electron mass

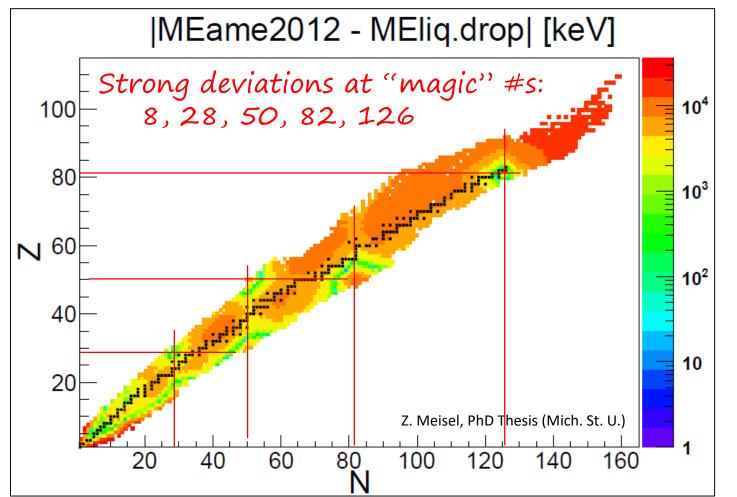
 $m_n = 939.56533 \text{ MeV/c}^2$ is the neutron mass

 $m_u = 931.494013 \text{ MeV/c}^2$ is the atomic mass unit

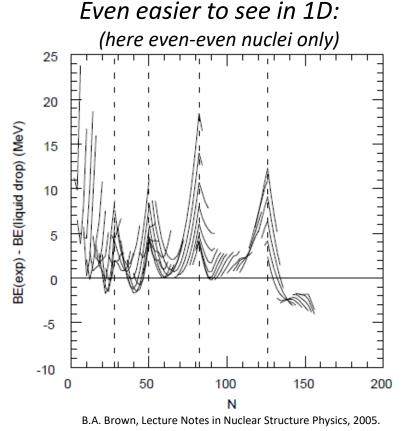
- Periodic evaluations are released for all known atomic masses, evaluated from a variety of experimental data by a group of noble souls.
 - This is referred to as the "Atomic Mass Evaluation" or "AME".
 - The latest is from 2016 (M. Wang et al. Ch. Phys. C 2017)

Binding Energy & Magic Numbers

 Binding energy per nucleon looks like a pretty smooth surface, what happens if we fit a smooth function to it & plot the residual? (we'll worry about what this smooth function is next time)



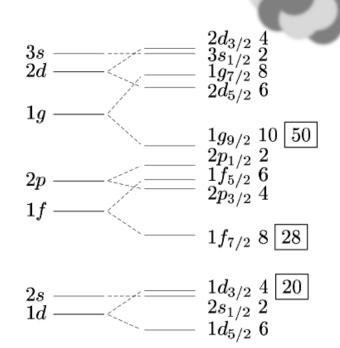




Magic Numbers and Nuclear Structure



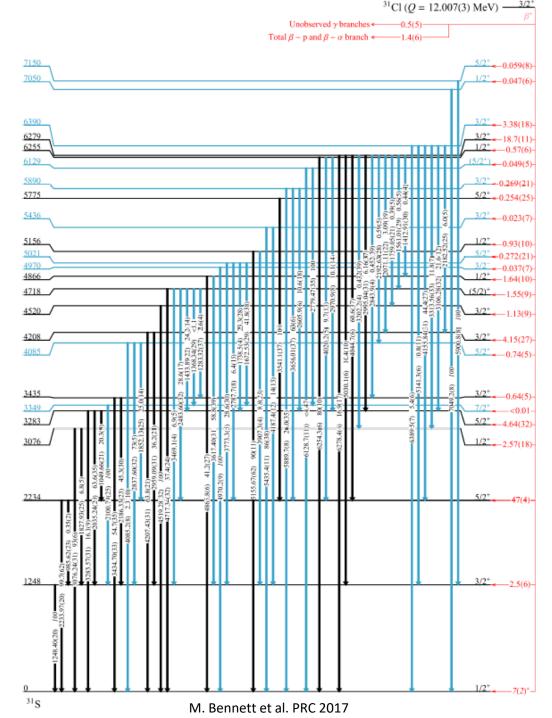
- Nucleons can only be combined in discrete configurations, signified as a "state"
- The lowest-energy state is the "ground state".
- A nucleus populated in a higher-lying state, an "excited state", will often γ -decay to de-excite to the ground state (though it may decay to a state of another nucleus).
- In the shell model (which we'll get to in later weeks), the properties of a state are calculated by considering neutrons and protons filling orbitals, like electrons for an atom.
- These are called single-particle levels:
- Note the gaps at the magic numbers. (We'll talk about how those come to be later)
- Exciting nucleons (or pairs) into other levels gives rise to excited states.



$$1s - - - 1s_{1/2} \ 2 \ 2$$

Nuclear Excited States

- The example on the right is for 31 S, where the levels shown were populated by β -decay from 31 Cl
- Note that the density of levels increases with excitation energy pretty dramatically.
- The only reason it doesn't keep increasing with excitation energy here is because, as we'll learn, this particular mechanism isn't suited for populating the higher lying states.
- Properties of the states include not only the excitation energy, but also
 - ullet Width: related to time to decay via γ or particle emission
 - Spin: intrinsic angular momentum
 - Parity: related to symmetry of wave function



Timeline of Selected Major Early Developments in Nuclear Physics

- 1896: Radioactivity discovered (uranium) (H. Becquerel, Comptes Rendus)
- 1905: Mass-energy equivalence proposed (A. Einstien, Annalen der Physik)
- 1911: Nuclear model of atom proposed (E. Rutherford, Phil. Mag.)
- 1913: Mass spectrometry invented (proton discovered) (J. Thomson, Proc. Phys. Soc. London)
- 1919: Isotope existence discovered (F. Aston, Nature)
- 1920: Nuclear transmutation proposed to power the sun (A. Eddington, The Observer)
- 1928: α-decay theory & barrier penetration (1st application of QM to real world problem) (G. Gamow, Z. Phys.)
- 1929: Liquid drop model of nucleus proposed (G. Gamow, Proc. Roy. Soc. A)
- 1932: Neutron discovered (J. Chadwick, Nature)
- 1932: Nuclei proposed to be interacting nucleons (based on neutron discovery) (W. Heisenberg, Z. Phys.)
- 1932: Nuclear transmutation using a particle accelerator (7 Li(p,α) α) (J. Cockroft & E. Walton, Proc. Phys. Soc. London)
- 1934: Theory of β decay (rejected from Nature & caused Fermi to switch to nuclear experiment) (E. Fermi, Z. Phys.)
- 1935: Semi-empirical mass formula developed (based on liquid drop) (C. Wiezsäcker, Z. Phys)
- 1936: Fusion theory (based on liquid drop) (N. Bohr, Nature)
- 1937: Fission proposed (based on liquid drop) (N. Bohr & J. Wheeler, Science)
- 1939: Fission measured (O. Hahn & F. Strassmann, Die Naturwissenschaften)
- 1939: Fission theory (L. Meitner & O. Frisch, Nature)
- 1949: Magic number explanation (M. Goeppert-Mayer, Phys. Rev.)
- 1952: Compound nuclear reaction formalism (W. Hauser & H. Feshbach, Phys. Rev.)
- 1957: Comprehensive theory of nucleosynthesis (M. Burbidge, G. Burbidge, W. Fowler, F. Hoyle, Rev Mod Phys.)
- 1957: Parity violation discovered (C. Wu et al., Phys. Rev)

Further Reading

- Chapters 1 & 2: Modern Nuclear Chemistry (Loveland, Morrissey, Seaborg)
- Chapters 1 & 2: Nuclear & Particle Physics (B.R. Martin)
- Chapter 1: Lecture Notes in Nuclear Structure Physics (B.A. Brown)